Exercise 12.1 *"Massive photon" in the Abelian Higgs model* (1.5 points)

Consider a model with a complex scalar field $\phi(x)$ whose dynamics is governed by the Lagrangian

$$\mathcal{L}_\phi(\phi, \partial\phi) = |\partial\phi|^2 - V(\phi^*\phi),$$

where V represents a general potential for the scalar self-interactions. Moreover, the quanta of ϕ carry the electric charge q, so that the full Lagrangian for ϕ including its electromagnetic interaction reads

$$\mathcal{L} = \mathcal{L}_\phi(\phi, D\phi) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

with the covariant derivative $D_\mu = \partial_\mu + iqA_\mu$ and the field-strength tensor $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ of the elmg. field $A^\mu(x)$.

a) The potential $V(\phi)$ is assumed to have a minimum for $|\phi(x)| \equiv v/\sqrt{2}$, so that this condition characterizes the ground state of the system (= field configuration of lowest energy). This suggests the following parametrization of ϕ:

$$\phi(x) = \frac{1}{\sqrt{2}}\Big(v + h(x)\Big) \exp\Big\{i\chi(x)/v\Big\},$$

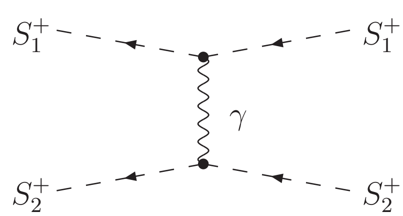
where $h(x)$ and $\chi(x)$ are real fields. Express $\mathcal{L}$ in terms of $h(x)$, $\chi(x)$, and $A^\mu(x)$.

b) Show that the model respects the gauge symmetry $\phi(x) \to \phi'(x) = \exp\{-iq\omega(x)\}\phi(x)$, $A_\mu(x) \to A'_\mu(x) = A_\mu(x) + \partial_\mu\omega(x)$ and argue that $\chi(x)$ is an unphysical field, i.e. that it can be consistently set to zero.

c) From $\mathcal{L}$ with $\chi = 0$, read off the part $\mathcal{L}_{AA,0}$ that is responsible for the free motion of A. Derive the Euler-Lagrange equation from $\mathcal{L}_{AA,0}$ for the free motion of A and identify the mass M_A of the quanta of A upon comparing the equation of motion with Proca's equation.

Exercise 12.2 *Elmg. scattering of two charged scalars* (2 points)

Consider the reaction $S_1^+(p_1) + S_2^+(p_2) \to S_1^+(k_1) + S_2^+(k_2)$ in scalar quantum electrodynamics, i.e. the particles $S_a^\mp$ $(a = 1, 2)$ are scalar particles with electric charges $\pm Q_a e$ and masses M_a. The corresponding Feynman rules are given on the back side. In Born approximation only the following diagram is relevant.



In the centre-of-mass frame the particle momenta read

$$p^\mu_{1,2} = (E_{1,2}, 0, 0, \pm p), \qquad k^\mu_{1,2} = (E_{1,2}, \pm p \sin\theta\cos\varphi, \pm p\sin\theta\sin\varphi, \pm p\cos\theta),$$

where $E_{1,2}$ are the energies of the incoming particles, and $p = \sqrt{E_a^2 - M_a^2}$ the absolute value of their three-momenta.

a) Calculate the squared transition matrix element $|\mathcal{M}|^2$ as function of the Mandelstam variables $s = (p_1 + p_2)^2$, $t = (p_1 - k_1)^2$, and $u = (p_1 - k_2)^2$.

b) Calculate the differential cross section

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s}|\mathcal{M}|^2 \,.$$

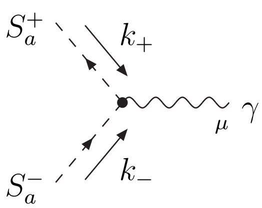
Compare the non-relativistic limit $(m_1, m_2 \gg p)$ of the result with the classical Rutherford cross section

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{Rutherford}} = \frac{\alpha^2 Q_1^2 Q_2^2}{4M^2\, v^4\, \sin^4(\frac{\theta}{2})},$$

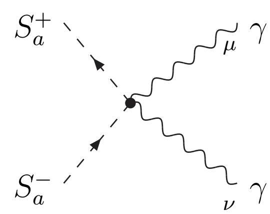
where in this limit $v = p/M$ and $M = M_1 M_2/(M_1 + M_2)$ denote the relative velocity and the reduced mass of the two-body system, respectively, and $\alpha = e^2/(4\pi)$ is the fine-structure constant.

<u>Feynman rules for the charged spin-0 particles $S_a^\pm$:</u>

Vertices:

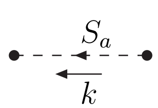


$iQ_a e(k_+ - k_-)_\mu$



$2iQ_a^2 e^2 g_{\mu\nu}$

Propagators and external lines:



$$\frac{i}{k^2 - M_a^2}$$

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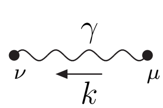

$$\frac{-i}{k^2}\left(g^{\mu\nu} - \frac{k^\mu k^\nu}{k^2}(1-\xi)\right)$$

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The fields $S_a^\pm$ are defined to be incoming at the vertices. Outgoing fields $S_a^\pm$ correspond to incoming fields $S_a^\mp$ with reversed momenta.